

## PROBLEM 4

4. Find the values of  $\mu$  where solutions bifurcate and examine the stability of the origin in each case. (Problems 4af on Page 108, PDF Page 131 )

a)  $x' = x + \mu y, y' = \mu x + y$

f)  $x' = y, y' = x^2 - x + \mu y$

**solution**

a) The nullclines:  $0 = x + \mu y, 0 = \mu x + y$

Of course,  $(0, 0)$  is one critical point

$$A = \begin{pmatrix} 1 & \mu \\ \mu & 1 \end{pmatrix}, \quad |A - \lambda I| = \lambda^2 - p\lambda + q = \lambda^2 - 2\lambda + (1 - \mu^2) = 0$$

With **Vieta theorem**,  $\Delta = 2^2 - 4 \times (1 - \mu^2) = 4\mu^2, \lambda_1 + \lambda_2 = p = 2 > 0, \lambda_1\lambda_2 = q = 1 - \mu^2$

For different values of  $\mu$

**condition 1:**  $q = 0, \Delta \neq 0$ , one eigenvalue  $\lambda_1 = 0$

$$q = 1 - \mu^2, \Delta = 4\mu^2 \neq 0 \Leftrightarrow \mu = \pm 1$$

the other eigenvalue  $\lambda_2 = p = 2 > 0$

critical point is a **unstable borderline**

**condition 2:**  $\Delta = 0, q \neq 0$ , repeated real eigenvalues  $\lambda_1 = \lambda_2 = \frac{p}{2}$

$$\Delta = 4\mu^2 = 0, q = 1 - \mu^2 \neq 0 \Leftrightarrow \mu = 0$$

repeated real eigenvalues  $\lambda_1 = \lambda_2 = \frac{p}{2} = 1 > 0$

critical point is a **unstable node**

$(A - \lambda I)w = 0$  eigenvector  $w_1, w_2$ : **proper node: star**  $X = (c_1w_1 + c_2w_2)e^{\lambda t}$

**condition 3:**  $\Delta > 0, q > 0$ , eigenvalues are both positive  $\lambda_1 > 0, \lambda_2 > 0$

$$\Delta = 4\mu^2 > 0, q = 1 - \mu^2 > 0 \Leftrightarrow -1 < \mu < 1, \mu \neq 0$$

critical point is an **unstable node**

**condition 4:**  $\Delta > 0, q < 0$ , signs fo eigenvalues are opposite  $\lambda_1 > 0, \lambda_2 < 0$

$$\Delta = 4\mu^2 > 0, q = 1 - \mu^2 < 0 \Leftrightarrow \mu < -1 \text{ or } 1 < \mu$$

critical point is a **unstable saddle**

f) The nullclines:  $0 = y, 0 = x^2 - x + \mu y$ , of course,  $(0, 0)$  is one critical point  
Consider the Jacobian  $J$ , at the critical point  $(0, 0)$

$$J = \begin{pmatrix} \frac{\partial P}{\partial x} & \frac{\partial P}{\partial y} \\ \frac{\partial Q}{\partial x} & \frac{\partial Q}{\partial y} \end{pmatrix} = \begin{pmatrix} 0 & 1 \\ 2x - 1 & \mu \end{pmatrix}, \quad J_{(0,0)} = \begin{pmatrix} 0 & 1 \\ -1 & \mu \end{pmatrix}$$

$$|J - \lambda I| = \lambda^2 - p\lambda + q = \lambda^2 - \mu\lambda + 1 = 0$$

With **Vieta theorem**,  $\Delta = \mu^2 - 4 \times 1 = \mu^2 - 4$ ,  $\lambda_1 + \lambda_2 = p = \mu > 0$ ,  $\lambda_1\lambda_2 = q = 1$

**condition 1:**  $\Delta = 0, q \neq 0$ , repeated real eigenvalues  $\lambda_1 = \lambda_2 = \frac{p}{2}$

$\Delta = \mu^2 - 4 = 0, q = 1 \neq 0 \Leftrightarrow \mu = \pm 2$

- $\mu = -2$  repeated real eigenvalues  $\lambda_1 = \lambda_2 = \frac{p}{2} = \frac{\mu}{2} = -1 < 0$

critical point is a **stable node**

$(J - \lambda I)w = 0, (J - \lambda I)v = w$ : **improper node**  $X = c_1 w e^{\lambda t} + c_2(w + vt)e^{\lambda t}$

- $\mu = +2$  repeated real eigenvalues  $\lambda_1 = \lambda_2 = \frac{p}{2} = \frac{\mu}{2} = +1 > 0$

critical point is an **unstable node**

$(J - \lambda I)w = 0, (J - \lambda I)v = w$ : **improper node**  $X = c_1 w e^{\lambda t} + c_2(w + vt)e^{\lambda t}$

**condition 2:**  $\Delta > 0, q > 0$ , eigenvalues have the same signs  $\lambda_1 < \lambda_2 < 0$  or  $\lambda_1 > \lambda_2 > 0$

$\Delta = \mu^2 - 4 > 0, q = 1 > 0 \Leftrightarrow \mu < -2$  or  $2 < \mu$

- $p = \mu < 0 \Rightarrow \mu < -2$  real eigenvalues  $\lambda_1 < \lambda_2 < 0$

critical point is a **stable node**

- $p = \mu > 0 \Rightarrow 2 < \mu$  real eigenvalues  $\lambda_1 > \lambda_2 > 0$

critical point is an **unstable node**

**condition 3:**  $\Delta < 0, p = 0$ , real part of eigenvalues  $\text{Re}(\lambda_1) = \text{Re}(\lambda_2) = \frac{p}{2} = 0$

$\Delta = \mu^2 - 4 < 0, p = \mu = 0 \Leftrightarrow \mu = 0$

we can't determine if the **spiral** is stable with Jacobian

solve  $x' = y, y' = x^2 - x \Rightarrow x'' = x^2 - x$ , with the substitution  $x'' = \frac{dx'}{dx} \frac{dx}{dt} = \frac{dy}{dx} y$

$$\frac{dy}{dx} y = x^2 - x \Leftrightarrow \int y dy = \int (x^2 - x) dx \Leftrightarrow x^2 + y^2 - \frac{2}{3} x^3 = C$$

critical point is a **stable center**

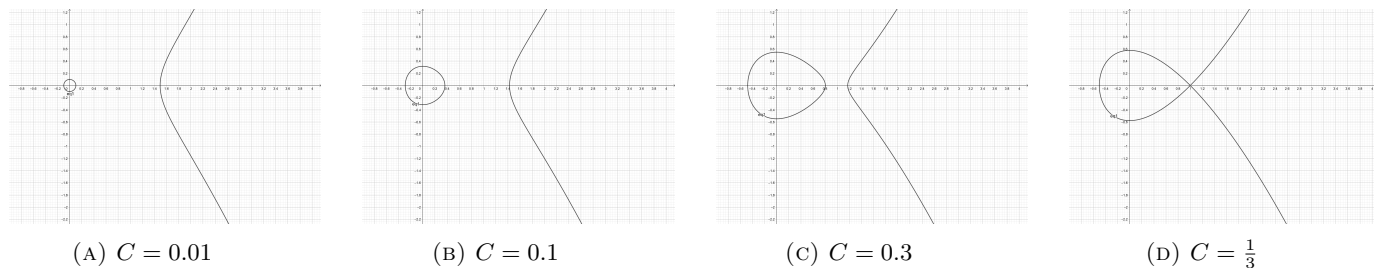


FIGURE 1. the diagram of  $p, \frac{dp}{d\tau}$

**condition 4:**  $\Delta < 0, p > 0$ , real part of eigenvalues are both positive  $\text{Re}(\lambda_1) = \text{Re}(\lambda_2) = \frac{p}{2} > 0$

$\Delta = \mu^2 - 4 < 0, p = \mu > 0 \Leftrightarrow 0 < \mu < 2$

critical point is an **unstable spiral**

**condition 5:**  $\Delta < 0, p > 0$ , real part of eigenvalues are both negative  $\text{Re}(\lambda_1) = \text{Re}(\lambda_2) = \frac{p}{2} < 0$

$\Delta = \mu^2 - 4 < 0, p = \mu < 0 \Leftrightarrow -2 < \mu < 0$

critical point is a **stable spiral**

## PROBLEM 10

10. (Ecology) Let  $P$  denote the carbon biomass of plants in an ecosystem,  $H$  be the carbon biomass of herbivores, and  $\phi$  the rate of primary carbon production in plants due to photosynthesis. A model of plant-herbivore dynamics is given by (Problem 10 on Page 109, PDF Page 133)

$$P' = \phi - aP - bHP$$

$$H' = \varepsilon bHP - cH$$

where  $a, b, c$ , and  $\varepsilon$  are positive parameters.

- Explain the various terms and parameters in the model and determine the dimensions of each parameter.
- Non-dimensionalize the model and find the equilibrium solutions.
- Analyze the dynamics in the cases of high primary production ( $\phi > ac/\varepsilon b$ ) and low primary production ( $\phi < ac/\varepsilon b$ ). Explain what happens in the system if primary production is slowly increased from a low value to a high value.

**solution**

a) For the various terms,  $[P] = [H] = M$ ,  $[\phi] = MT^{-1}$

For the parameters  $a, b, c$ ,  $[aP] = [a]M = MT^{-1}$ ,  $[bHP] = [b]M^2 = MT^{-1}$ ,  $[cH] = [c]M = MT^{-1}$ , thus  $[a] = T^{-1}$ ,  $[b] = M^{-1}T^{-1}$ ,  $[c] = T^{-1}$

For the parameter  $\varepsilon$ ,  $[\varepsilon bHP] = [\varepsilon]M^{-1}T^{-1}M^2 = MT^{-1}$ , thus  $[\varepsilon] = 1$  is dimensionless

b) Non-dimensionalize the model

Firstly, the characteristic scale  $t_c$  to make the coefficient of  $cH$  becomes 1

$$\frac{1}{t_c} = c \Rightarrow t_c \equiv \frac{1}{c}$$

Secondly, the  $P_c$  to make the coefficient of  $\varepsilon bHP$  becomes 1

$$\frac{1}{t_c} = \varepsilon b P_c \Rightarrow P_c \equiv \frac{c}{\varepsilon b}$$

Next, the  $H_c$  to make the coefficient of  $bHP$  becomes 1

$$\frac{1}{t_c} = b H_c \Rightarrow H_c \equiv \frac{c}{b}$$

In the end, with the substitution  $\tau \equiv t/t_c$ ,  $p = P/P_c$ ,  $h \equiv H/H_c$

$$\frac{dp}{d\tau} = -\left(\frac{a}{c}\right)\left(p - \frac{\phi}{[ac/\varepsilon b]}\right) - hp$$

$$\frac{dh}{d\tau} = hp - h$$

Define parameters  $\gamma \equiv \frac{a}{c}$ ,  $\phi_{th} \equiv [ac/\varepsilon b]$ ,  $\varphi \equiv \frac{\phi}{\phi_{th}}$

$$\frac{dp}{d\tau} = -\gamma\left(p - \frac{\phi}{\phi_{th}}\right) - hp = -\gamma(p - \varphi) - hp$$

$$\frac{dh}{d\tau} = hp - h$$

(1)

c) The nullclines of  $h$ :  $0 = -\gamma(p - \varphi) - hp \Leftrightarrow h = \gamma(-1 + \frac{\varphi}{p})$

the nullclines of  $p$ :  $0 = hp - h \Leftrightarrow h = 0, p = 1$

Consider the Jacobian  $J$

$$J = \begin{pmatrix} \frac{\partial P}{\partial x} & \frac{\partial P}{\partial y} \\ \frac{\partial Q}{\partial x} & \frac{\partial Q}{\partial y} \end{pmatrix} = \begin{pmatrix} -\gamma - h & -p \\ h & p - 1 \end{pmatrix}$$

**condition 1**  $\varphi \equiv \frac{\phi}{\phi_{th}} \equiv \frac{\phi}{[ac/\varepsilon b]} = 1$

the critical point  $(1, 0)$  for all  $p \geq 0, h \geq 0$

$$J = \begin{pmatrix} -\gamma & -1 \\ 0 & 0 \end{pmatrix}, \quad \det(J) = 0, |J - \lambda I| = \lambda^2 + \gamma\lambda = 0$$

eigenvalues  $\lambda_1 = 0, \lambda_2 = -\gamma < 0$  the equilibria  $(1, 0)$  has the degenerated type: **concentrated in a line**  $-\gamma(p - 1) - h = 0$  (the line go through critical point), is a **stable borderline**

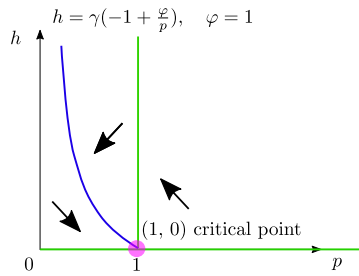


FIGURE 2. the phase diagram of  $\varphi = 1$

**condition 2**  $0 < \varphi \equiv \frac{\phi}{\phi_{th}} \equiv \frac{\phi}{[ac/\varepsilon b]} < 1$

the critical point  $(\varphi, 0)$  for all  $p \geq 0, h \geq 0$

$$J = \begin{pmatrix} -\gamma & -\varphi \\ 0 & \varphi - 1 \end{pmatrix}, \quad |J - \lambda I| = [\lambda - (-\gamma)][\lambda - (\varphi - 1)] = 0$$

eigenvalues  $\lambda_1 = \varphi - 1 < 0, \lambda_2 = -\gamma < 0$  the critical point  $(\varphi, 0)$  is a **stable node**

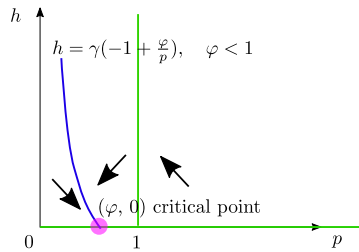


FIGURE 3. the phase diagram of  $0 < \varphi < 1$

**condition 3**  $1 < \varphi \equiv \frac{\phi}{\phi_{th}} \equiv \frac{\phi}{[ac/\varepsilon b]}$

the critical points  $(\varphi, 0), (1, \gamma(\varphi - 1))$  for all  $p \geq 0, h \geq 0$

- At  $(\varphi, 0)$

$$J = \begin{pmatrix} -\gamma & -\varphi \\ 0 & \varphi - 1 \end{pmatrix}, \quad |J - \lambda I| = [\lambda - (-\gamma)][\lambda - (\varphi - 1)] = 0$$

eigenvalues  $\lambda_1 = \varphi - 1 > 0, \lambda_2 = -\gamma < 0$  the critical point  $(\varphi, 0)$  is a **unstable saddle**

- At  $(1, \gamma(\varphi - 1))$

$$J = \begin{pmatrix} -\gamma\varphi & -1 \\ \gamma(\varphi - 1) & 0 \end{pmatrix}, \quad |J - \lambda I| = \lambda^2 + \gamma\varphi\lambda + \gamma(\varphi - 1) = \lambda^2 - p\lambda + q = 0$$

With **Vieta theorem**  $p = \lambda_1 + \lambda_2 = -\gamma\varphi < 0$ ,  $q = \lambda_1\lambda_2 = \gamma(\varphi - 1) > 0$ ,  $\Delta = (\gamma\varphi)^2 - 4\gamma(\varphi - 1) = (\gamma\varphi - 2)^2 + 4(\gamma - 1)$

the critical point  $(1, \gamma(\varphi - 1))$  is always **stable**

(1)  $\gamma \equiv \frac{a}{c} = 1$

(a) if  $\varphi = \frac{2}{\gamma} = 2$

the repeated eigenvalues  $\lambda_1 = \lambda_2 = -1$ , like  $c_1we^{\lambda_1 t} + c_2(w + vt)e^{\lambda_1 t}$

the critical point  $(1, \gamma(\varphi - 1)) = (1, 1)$  is a **stable improper node**

(b) if  $\varphi \neq \frac{2}{\gamma} = 2$

hence,  $\Delta > 0$  always holds, 2 real eigenvalues  $\lambda_1 < \lambda_2 < 0$

the critical point  $(1, \gamma(\varphi - 1))$  is a **stable node**

(2)  $\gamma \equiv \frac{a}{c} > 1$

hence,  $\Delta > 0$  always holds, it has 2 real eigenvalues  $\lambda_1 < \lambda_2 < 0$

the critical point  $(1, \gamma(\varphi - 1))$  is a **stable node**

(3)  $0 < \gamma \equiv \frac{a}{c} < 1$

(a) if  $\varphi = \frac{2(1 \pm \sqrt{1 - \gamma})}{\gamma}$

hence  $\Delta = 0$ , the repeated eigenvalues  $\lambda_1 = \lambda_2 = -1 \mp \sqrt{1 - \gamma}$ , solution is in form of  $c_1we^{\lambda_1 t} + c_2(w + vt)e^{\lambda_1 t}$

the critical point  $(1, \gamma(\varphi - 1)) = (1, 2 - \gamma \pm 2\sqrt{1 - \gamma})$  is a **stable improper node**

(b) if  $\frac{2(1 - \sqrt{1 - \gamma})}{\gamma} < \varphi < \frac{2(1 + \sqrt{1 - \gamma})}{\gamma}$

hence  $\Delta < 0$ , the complex eigenvalues  $\text{Re}(\lambda_1) = \text{Re}(\lambda_2) = \frac{p}{2} = -\gamma\varphi/2 < 0$

the critical point  $(1, \gamma(\varphi - 1))$  is a **stable spiral**

(c) if  $0 < \varphi < \frac{2(1 - \sqrt{1 - \gamma})}{\gamma}$  or  $\frac{2(1 + \sqrt{1 - \gamma})}{\gamma} < \varphi$

hence  $\Delta > 0$ , 2 real eigenvalues  $\lambda_1 < \lambda_2 < 0$

the critical point  $(1, \gamma(\varphi - 1))$  is a **stable node**

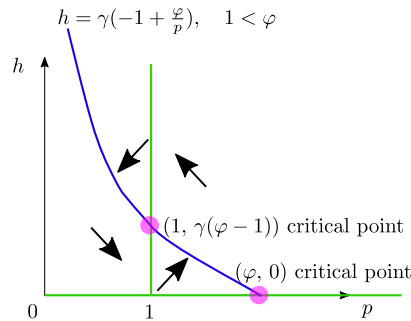


FIGURE 4. the phase diagram of  $1 < \varphi$

As  $\phi$  increases from low to high, the equilibrium of the biomass  $(P^*, H^*)$  changes as follows:

- For  $0 < \varphi \equiv \frac{\phi}{[ac/\epsilon b]} \leq 1$

**stable** critical point  $(p, h) = (\varphi, 0) \Leftrightarrow$  equilibria  $(P^*, H^*) = (\varphi P_c, 0) = (\frac{\phi}{[ac/\epsilon b]} \frac{c}{\epsilon b}, 0) = (\frac{\phi}{a}, 0)$

- For  $1 < \varphi \equiv \frac{\phi}{[ac/\epsilon b]}$

**stable** critical point  $(p, h) = (1, \gamma(\varphi - 1))$

$\Leftrightarrow$  equilibria  $(P^*, H^*) = (P_c, \gamma(\varphi - 1)H_c) = (\frac{c}{\epsilon b}, \frac{a}{c}(\frac{\phi}{[ac/\epsilon b]} - 1)\frac{c}{b}) = (\frac{c}{\epsilon b}, \frac{\epsilon\phi}{c} - \frac{a}{b})$

## PROBLEM 9

To find approximations to the roots of the cubic equation

$$x^3 - 4.001x + 0.002 = 0$$

why is it easier to examine the equation (Problem 9 on Page 166, PDF Page 199)

$$x^3 - (4 + \varepsilon)x + 2\varepsilon = 0$$

- Find a two-term approximation to this equation.
- Then use software (e.g. matlab) to solve the above equation and compare the results.

**solution**

**Perturbation** is successful when  $x_{approx.}(\varepsilon) - x(\varepsilon) \rightarrow 0$  at some well-defined rate as  $\varepsilon \rightarrow 0$

We can find 3 solutions for  $x^3 - (4 + \varepsilon)x + 2\varepsilon = 0$  when  $\varepsilon = 0$ , that is  $x_0 = 0, -2, +2$ , corresponding to 3 continuous functions  $x = f_1(\varepsilon), f_2(\varepsilon), f_3(\varepsilon)$ , and  $f_1(0) = 0, f_2(0) = -2, f_3(0) = +2$

Expand  $x = f_i(\varepsilon), i = 1, 2, 3$

$$f_i(\varepsilon) = x_0 + x_1\varepsilon + x_2\varepsilon^2 + \dots$$

Substitute  $x = f_i(x) = x_0 + x_1\varepsilon + x_2\varepsilon^2 + \dots$  in  $x^3 - (4 + \varepsilon)x + 2\varepsilon = 0$

Compare the coefficients of  $1, \varepsilon, \varepsilon^2$  respectively

$$1 : x_0^3 - 4x_0 = 0$$

$$\varepsilon : \binom{3}{2} x_0^2 x_1 - 4x_1 - x_0 + 2 = 0$$

$$\varepsilon^2 : \binom{3}{1} x_0 x_1^2 + \binom{3}{2} x_0^2 x_2 - 4x_2 - x_1 = 0$$

For  $f_1(\varepsilon), f_1(0) = x_0 = 0$ , we can conclude consequently

$$x_0 = 0 \Rightarrow -4x_1 + 2 = 0 \Rightarrow x_1 = \frac{1}{2} \Rightarrow -4x_2 - \frac{1}{2} = 0 \Rightarrow x_2 = -\frac{1}{8}$$

$$f_1(\varepsilon) = 0 + \frac{1}{2}\varepsilon - \frac{1}{8}\varepsilon^2 + \dots, \quad x = f_1(0.001) \approx 0 + \frac{1}{2}(0.001) - \frac{1}{8}(0.001)^2 = 0.000499875$$

For  $f_2(\varepsilon), f_2(0) = x_0 = -2$ , we can conclude consequently

$$x_0 = -2 \Rightarrow 12x_1 - 4x_1 + 2 + 2 = 0 \Rightarrow x_1 = -\frac{1}{2} \Rightarrow -\frac{3}{2} + 12x_2 - 4x_2 + \frac{1}{2} = 0 \Rightarrow x_2 = \frac{1}{8}$$

$$f_2(\varepsilon) = -2 - \frac{1}{2}\varepsilon + \frac{1}{8}\varepsilon^2 + \dots, \quad x = f_2(0.001) \approx -2 - \frac{1}{2}(0.001) + \frac{1}{8}(0.001)^2 = -2.000499875$$

For  $f_3(\varepsilon), f_3(0) = x_0 = +2$ , we can conclude consequently

$$x_0 = +2 \Rightarrow 12x_1 - 4x_1 - 2 + 2 = 0 \Rightarrow x_1 = 0 \Rightarrow 0 + 12x_2 - 4x_2 - 0 = 0 \Rightarrow x_2 = 0 \Rightarrow x_n = 0 (n \geq 3)$$

$$f_3(\varepsilon) = 2 + 0\varepsilon + 0\varepsilon^2 + \dots = 2, \quad x = f_3(0.001) = 2$$

Finally we have 2 approximated roots  $x \approx 0.000499875, -2.000499875$  and one exact root  $x = 2$

```

1 >> syms x;
2 >> vpasolve(x^3 - 4.001*x + 0.002)
3 ans =
4 -2.0004998750624609648232582877001
5 0.00049987506246096482325828770010975
6 2.0

```

## PROBLEM 20

20. Find a two-term perturbation solution of

$$u' + u = \frac{1}{1 + \varepsilon u}, \quad u(0) = 0, \quad 0 < \varepsilon \ll 1$$

Use software to plot the approximate solution with  $\varepsilon = 0.1$  for  $0 \leq x \leq 1$ .

(Problem 20 on Page 169, PDF Page 201)

**solution**

We can find function  $u = u(\varepsilon, x)$ , and  $u(0, x)$  is solution for  $u' + u = \frac{1}{1+0} = 1, u(0) = 0$

Expand  $u(\varepsilon, x)$  with  $\varepsilon$

$$u(\varepsilon, x) = y_0(x) + y_1(x)\varepsilon + y_2(x)\varepsilon^2 + \dots$$

Substitute  $u(\varepsilon, x) = y_0(x) + y_1(x)\varepsilon + y_2(x)\varepsilon^2 + \dots$  in  $u' + u = \frac{1}{1+\varepsilon u}$

Compare the coefficients of  $1, \varepsilon, \varepsilon^2$  respectively

$$1 : y_0' + y_0 = 1$$

$$\varepsilon : y_1' + y_1 = -y_0$$

$$\varepsilon^2 : y_2' + y_2 = -y_1 + y_0$$

Initial condition  $u(\varepsilon, 0) = y_0(0) + y_1(0)\varepsilon + y_2(0)\varepsilon^2 + \dots = 0$ , compare the coefficients of  $\varepsilon^k, k \in \mathbb{N}$

$$y_0(0) = y_1(0) = y_2(0) = \dots = 0$$

Solve  $y_0$

$$y_0 = e^{-\int 1 dx} \left[ \int 1 \cdot e^{\int 1 dx} dx \right] = 1 + c_0 e^{-x}, y_0(0) = 0 \Rightarrow y_0 = 1 - e^{-x}$$

Solve  $y_1$

$$y_1 = e^{-x} \left[ \int -(1 - e^{-x}) \cdot e^x dx \right] = e^{-x} [-e^x + x + c_2] = -1 + x e^{-x} + c_1 e^{-x}, y_1(0) = 0 \Rightarrow y_1 = -1 + (x+1)e^{-x}$$

Solve  $y_2$

$$\begin{aligned} y_2 &= e^{-x} \left[ \int [ -(-1 + (x+1)e^{-x}) + 1 - e^{-x} ] \cdot e^x dx \right] = e^{-x} \left[ \int [2 - (x+2)e^{-x}] \cdot e^x dx \right] \\ &= e^{-x} \left[ 2e^x - \left( \frac{1}{2}x^2 + 2x + c_2 \right) \right], y_2(0) = 0 \\ \Rightarrow y_2 &= e^{-x} \left[ 2e^x - \left( \frac{1}{2}x^2 + 2x + 2 \right) \right] = 2 - \left( \frac{1}{2}x^2 + 2x + 2 \right) e^{-x} \end{aligned}$$

In all

$$y_0 = 1 - e^{-x}$$

$$y_1 = -1 + (x+1)e^{-x}$$

$$y_2 = 2 - \left( \frac{1}{2}x^2 + 2x + 2 \right) e^{-x}$$

$$u(\varepsilon, x) \approx y_0(x) + y_1(x)\varepsilon + y_2(x)\varepsilon^2$$

$$= [1 - e^{-x}] + \varepsilon [-1 + (x+1)e^{-x}] + \varepsilon^2 \left[ 2 - \left( \frac{1}{2}x^2 + 2x + 2 \right) e^{-x} \right]$$

The close form solution when  $\varepsilon = 0.1$  for the separable equation,  $C$  is given by  $u(0) = 0$

$$C - x = \frac{1}{14} \left[ (7 + \sqrt{35}) \ln | -u + \sqrt{35} - 5 | - (\sqrt{35} - 7) \ln | u + \sqrt{35} + 5 | \right]$$

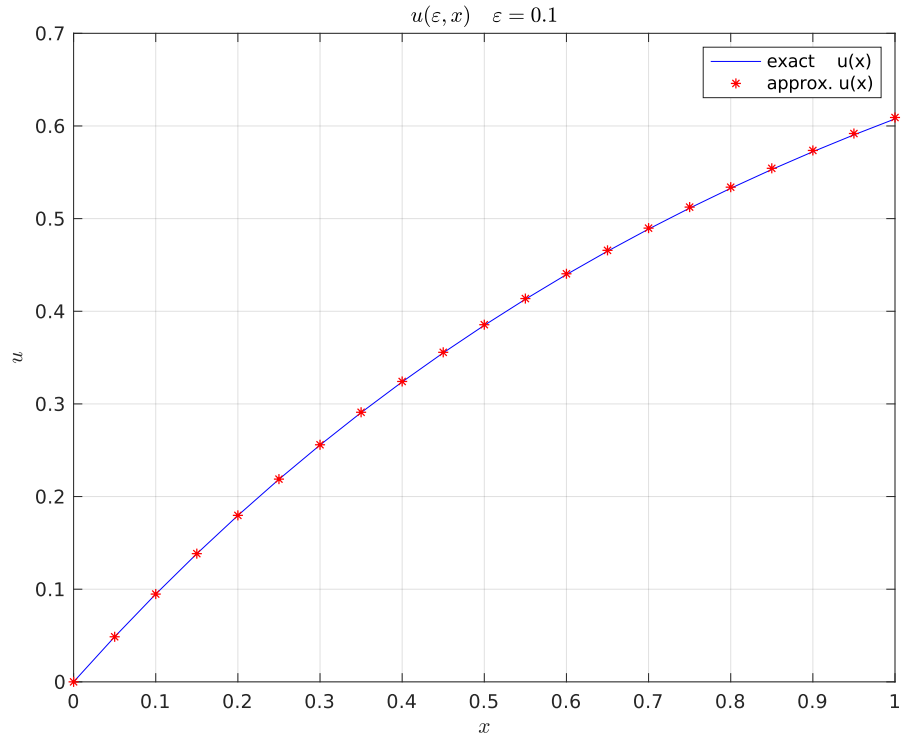


FIGURE 5. the exact  $u(\varepsilon, x)$  and the approximate  $u(\varepsilon, x)$  ( $\varepsilon = 0.1$ )

```

1 clear; clc; close all
2 % solve u(x) with ode23()
3 eps = 0.1;
4 func = @(x, u) 1 / (1 + eps*u) - u;
5 u0 = 0;
6 tspan = (0:0.05:1)';
7 [x,u] = ode23(func, tspan, u0);
8 plot(x,u, 'b-'); xlabel('$x$', 'Interpreter', 'latex');
9 ylabel('$u$', 'Interpreter', 'latex');
10 title('$u(\varepsilon, x)\quad\varepsilon=0.1$', 'Interpreter', 'latex');
11 % calc approximate u(x)
12 y0 = @(x) 1 - exp(-x);
13 y1 = @(x) -1 + (x + 1).* exp(-x);
14 y2 = @(x) 2 - (x.^2 / 2 + 2 * x + 2).* exp(-x);
15 u_perturb = @(x) y0(x) + eps * y1(x) + eps^2 * y2(x);
16 u_approx = u_perturb(tspan);
17 hold on; plot(tspan, u_approx, 'r*'); grid on;
18 legend('exact u(x)', 'approx. u(x)');

```



## BONUS.

16. (Circuits) An RCL circuit with a nonlinear resistance, where the voltage drop across the resistor is a nonlinear function of current, can be modeled by the Van der Pol equation

$$x'' + \rho(x^2 - 1)x' + x = 0$$

where  $\rho$  is a positive constant, and  $x(t)$  is the current.

- In the phase plane, show that the origin is an unstable equilibrium.
- Sketch the nullclines and the vector field. What are the possible dynamics? Is there a limit cycle?

(Problem 16ab on Page 111, PDF Page 134)

**solution**

With the substitution  $y = x'$

$$x' = y$$

$$y' = x'' = -x - \rho(x^2 - 1)y$$

The nullclines:  $0 = y, 0 = -x - \rho(x^2 - 1)y$   
the only critical point is  $(0, 0)$

For the direction in the separated regions:

**region 1:**  $y > 0, -x - \rho(x^2 - 1)y > 0$

$(x', y') = (y, -x - \rho(x^2 - 1)y) = (+, +)$

**region 2:**  $y < 0, -x - \rho(x^2 - 1)y > 0$

$(x', y') = (y, -x - \rho(x^2 - 1)y) = (-, +)$

**region 3:**  $y < 0, -x - \rho(x^2 - 1)y < 0$

$(x', y') = (y, -x - \rho(x^2 - 1)y) = (-, -)$

**region 4:**  $y > 0, -x - \rho(x^2 - 1)y < 0$

$(x', y') = (y, -x - \rho(x^2 - 1)y) = (+, -)$

Consider the Jacobian  $J$

$$J = \begin{pmatrix} \frac{\partial P}{\partial x} & \frac{\partial P}{\partial y} \\ \frac{\partial Q}{\partial x} & \frac{\partial Q}{\partial y} \end{pmatrix} = \begin{pmatrix} 0 & 1 \\ -1 - 2\rho xy & -\rho(x^2 - 1) \end{pmatrix}$$

a) At the critical point  $(0, 0)$ ,

$$J = \begin{pmatrix} 0 & 1 \\ -1 & \rho \end{pmatrix}, \quad \det(J) \neq 0, \quad |J - \lambda I| = \lambda^2 - \rho\lambda + 1 = \lambda^2 - p\lambda + q = 0$$

With **Vieta theorem**,  $\Delta = \rho^2 - 4, p = \lambda_1 + \lambda_2 = \text{Re}(\lambda_1) + \text{Re}(\lambda_2) = \rho > 0, q = 1$   
if stable  $\Rightarrow \text{Re}(\lambda_1) \leq 0, \text{Re}(\lambda_2) \leq 0 \Rightarrow \text{Re}(\lambda_1) + \text{Re}(\lambda_2) \leq 0$ , there is a conflict,  $(0, 0)$  is **unstable**

b) The the nullclines and the vector field are displayed in the figure above

(1)  $\Delta = \rho^2 - 4 = 0 \Rightarrow \rho = 2$

repeated eigenvalues  $\lambda_1 = \lambda_2 = 1$ , solution  $c_1 w e^{\lambda_1 t} + c_2(w + vt)e^{\lambda_1 t}$

the critical point  $(0, 0)$  is an **unstable improper node**

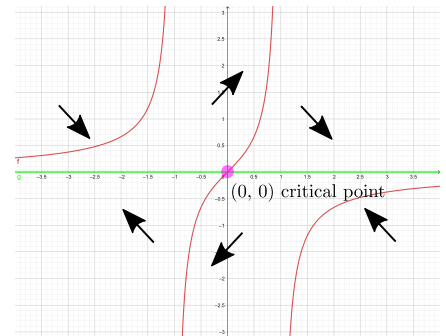


FIGURE 6. phase diagram  $\rho = 1$

- (2)  $\Delta = \rho^2 - 4 > 0 \Rightarrow 2 < \rho$   
 real eigenvalues  $\lambda_1 > \lambda_2 > 0$   
 the critical point  $(0, 0)$  is an **unstable node**
- (3)  $\Delta = \rho^2 - 4 < 0 \Rightarrow 0 < \rho < 2$   
 complex eigenvalues  $\text{Re}(\lambda_1) = \text{Re}(\lambda_2) = \frac{\rho}{2} = \frac{\rho}{2} > 0$   
 the critical point  $(0, 0)$  is an **unstable spiral**

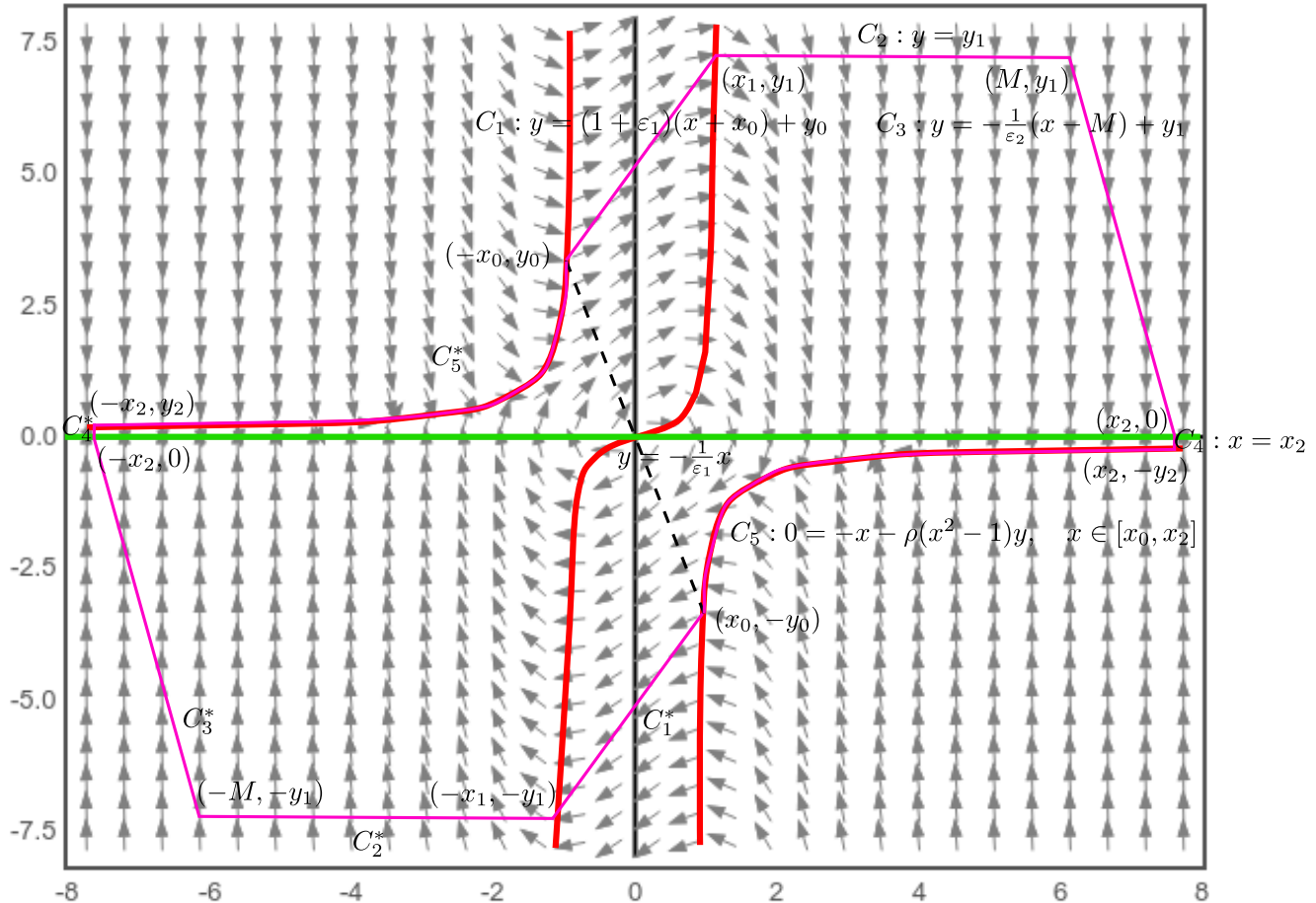


FIGURE 7. the simple closed curve  $C$ , where  $(x', y')^T \cdot \vec{n} < 0$

(Note: **Green**: x nullcline; **Red**: y nullcline; **Magenta**: the simple closed curve  $C$ )

**Theorem (Poincare-Bendixson Ring Domain Theorem).** Suppose  $R$  is the finite region of the plane lying between two simple closed curves  $C$  and  $\bar{C}$ , and  $F$  is the velocity vector field for the system  $x' = f(x, y)$   $y' = g(x, y)$ . If

- (i) at each point of  $C$  and  $\bar{C}$ , the field  $F$  points toward the interior of  $R$ , and  
 (ii)  $R$  contains no critical points,

then the system has a **closed trajectory** lying inside  $R$

See the [MIT limit cycle note](https://math.mit.edu/~jorloff/supnotes/supnotes03/lc.pdf) or go to the url: <https://math.mit.edu/~jorloff/supnotes/supnotes03/lc.pdf>

For **Poincare-Bendixson Ring Domain Theorem**, we construct the simple closed curves first. curve  $\mathbf{C}$  consists of  $C_1, C_2, C_3, C_4, C_5$  and the symmetrical curves  $C_1^*, C_2^*, C_3^*, C_4^*, C_5^*$  of  $(0, 0)$ . For  $(x, y)$  on  $\mathbf{C}$ , always holds  $(x', y')^T \cdot \vec{n} < 0$ , here  $(x', y')$ : velocity field,  $\vec{n}$ : normal vector at  $(x, y)$ . Namely, for all  $(x, y)$  on  $\mathbf{C}$ , velocity field  $(x', y')$  points toward the inside of  $\mathbf{C}$ . Short explanation for  $(x', y')^T \cdot \vec{n} < 0$  on  $C_1, C_2, C_3, C_4, C_5$  as follows:

(1)  $C_1 : y = (\rho + \varepsilon_1)(x + x_0) + y_0$

where  $(-x_0, y_0)$  is the intersection of  $y = -\frac{1}{\varepsilon_1}x$  and left branch of  $0 = -x - \rho(x^2 - 1)y$

$\varepsilon_1$  must satisfy, at  $(-x_0, y_0)$ :  $\frac{dy}{dx}|_{(-x_0, y_0)}$  slope of  $0 = -x - \rho(x^2 - 1)y > (\rho + \varepsilon_1)$  slope of  $C_1$

$$(-x_0, y_0) = \left( -\sqrt{1 + \frac{\varepsilon_1}{\rho}}, \sqrt{1 + \frac{\varepsilon_1}{\rho}}/\varepsilon_1 \right), \quad \text{left brach } 0 = -x - \rho(x^2 - 1)y \Leftrightarrow y = -\frac{1}{2\rho} \left( \frac{1}{x-1} + \frac{1}{x+1} \right)$$

$$\frac{dy}{dx} = \frac{1}{\rho} \frac{x^2 + 1}{(x^2 - 1)^2}, \quad \frac{dy}{dx}|_{(-x_0, y_0)} = \frac{1}{\rho} \frac{2 + \frac{\varepsilon_1}{\rho}}{(\frac{\varepsilon_1}{\rho})^2} = \frac{2\rho + \varepsilon_1}{\varepsilon_1^2} > (\rho + \varepsilon_1) \Leftrightarrow \rho > 0 > -\varepsilon_1 \left( \frac{1 - \varepsilon_1^2}{2 - \varepsilon_1^2} \right)$$

We can find  $0 < \varepsilon_1 < 1$  to satisfy it, for  $\forall \rho > 0$ , it always holds

$$(x', y')^T = (y, -x - \rho(x^2 - 1)y)^T, \quad \vec{n} = (-\rho - \varepsilon_1, 1)^T$$

$$(x', y')^T \cdot \vec{n} = -[(\varepsilon_1 y + x) + \rho x^2] < 0$$

(2)  $C_2 : y = y_1$

where  $(x_1, y_1)$  is the intersection of  $C_1$  and center branch of  $0 = -x - \rho(x^2 - 1)y$

for  $(x, y)$  on  $C_2$ , it has  $x > x_1 > 0, y = y_1 > 0$ , always holds

$$(x', y')^T = (y, -x - \rho(x^2 - 1)y)^T, \quad \vec{n} = (0, 1)^T$$

$$(x', y')^T \cdot \vec{n} = -x - \rho(x^2 - 1)y = [-x_1 - \rho(x_1^2 - 1)y_1] - (x - x_1) - (x^2 - x_1^2)y_1 = -(x - x_1) - (x^2 - x_1^2)y_1 < 0$$

(3)  $C_3 : y = -\frac{1}{\varepsilon_2}(x - M) + y_1$

where  $(M, y_1)$  satisfies  $M > \sqrt{1 + \frac{1}{\rho\varepsilon_2}}$

for  $(x, y)$  on  $C_3$ , it has  $x > M, y > 0$ , always holds

$$(x', y')^T = (y, -x - \rho(x^2 - 1)y)^T, \quad \vec{n} = (1, \varepsilon_2)^T$$

$$(x', y')^T \cdot \vec{n} = -\varepsilon_2 x - y [\rho\varepsilon_2(x^2 - 1) - 1] = -\varepsilon_2 x - y [\rho\varepsilon_2(M^2 - 1) - 1] - y\rho\varepsilon_2(x^2 - M^2) < 0$$

(4)  $C_4 : x = x_2$

where  $(x_2, 0)$  is the intersection of  $C_3$  and  $y = 0$

for  $(x, y)$  on  $C_4$ , it has  $x = x_2, y < 0$ , always holds

$$(x', y')^T = (y, -x - \rho(x^2 - 1)y)^T, \quad \vec{n} = (1, 0)^T$$

$$(x', y')^T \cdot \vec{n} = y < 0$$

(5)  $C_5 : 0 = -x - \rho(x^2 - 1)y$  right branch

where  $(x_2, -y_2)$  is the intersection of  $C_4$  and  $0 = -x - \rho(x^2 - 1)y$  right branch

for  $(x, y)$  on  $C_5$ , it has  $x > x_0 = \sqrt{1 + \frac{\varepsilon_1}{\rho}} > 1, y < 0$ , always holds

$$y = -\frac{1}{2\rho} \left( \frac{1}{x-1} + \frac{1}{x+1} \right), \quad \frac{dy}{dx} = \frac{1}{\rho} \frac{x^2 + 1}{(x^2 - 1)^2}$$

$$(x', y')^T = (y, -x - \rho(x^2 - 1)y)^T = (y, 0)^T, \quad \vec{n} = \left( \frac{dy}{dx}, -1 \right)^T$$

$$(x', y')^T \cdot \vec{n} = - \left[ \frac{1}{\rho} \frac{x^2 + 1}{(x^2 - 1)^2} \right] (-y) < 0$$

To sum up,  $(x', y')^T \cdot \vec{n} < 0$  holds on  $C_1, C_2, C_3, C_4, C_5$ , the velocity field is symmetrical to  $(0, 0)$ . For the symmetrical curves  $C_1^*, C_2^*, C_3^*, C_4^*, C_5^*$  of  $(0, 0)$ ,  $(x', y')^T \cdot \vec{n} < 0$  holds on  $C_1^*, C_2^*, C_3^*, C_4^*, C_5^*$ . Thus,  $(x', y')^T \cdot \vec{n} < 0$  holds for almost all points on entire  $\mathbf{C}$

At each point of  $\mathbf{C}$ , the field  $\mathbf{F} = (x', y')^T$  points toward the inside of  $\mathbf{C}$

Now, consider to construct the simple closed curve  $\bar{\mathbf{C}}$  (Magenta) near the critical point  $(0, 0)$

(1)  $\rho > 2, \Delta < 0$ , construct  $\bar{\mathbf{C}}$  as below

eigenvalues  $\lambda_{1,2} = \frac{p \pm \sqrt{p^2 - 4}}{2} > 0$ , with  $(J - \lambda_{1,2})w_{1,2} = 0$ , find  $w_{1,2} = [1, \frac{p \pm \sqrt{p^2 - 4}}{2}]^T$   
then  $X(t) = (x, y)^T = c_1 e^{\lambda_1 t} w_1 + c_2 e^{\lambda_2 t} w_2$ ,  $(x', y')^T = c_1 \lambda_1 e^{\lambda_1 t} w_1 + c_2 \lambda_2 e^{\lambda_2 t} w_2$   
notice that coefficient  $(c_1 \lambda_1 e^{\lambda_1 t}, c_2 \lambda_2 e^{\lambda_2 t})$  has the same signs as  $(c_1 e^{\lambda_1 t}, c_2 e^{\lambda_2 t})$ ,  $(c_1, c_2)$   
the normal vector at  $(x, y)$  is  $\vec{n} = \text{sgn}(c_1) \frac{w_1}{|w_1|} + \text{sgn}(c_2) \frac{w_2}{|w_2|}$ , it always holds

$$\begin{aligned} (x', y')^T \cdot \vec{n} &= (c_1 \lambda_1 e^{\lambda_1 t} w_1 + c_2 \lambda_2 e^{\lambda_2 t} w_2) \cdot \left( \text{sgn}(c_1) \frac{w_1}{|w_1|} + \text{sgn}(c_2) \frac{w_2}{|w_2|} \right) \\ &= \left( c_1 \lambda_1 e^{\lambda_1 t} |w_1| \frac{w_1}{|w_1|} + c_2 \lambda_2 e^{\lambda_2 t} |w_2| \frac{w_2}{|w_2|} \right) \cdot \left( \text{sgn}(c_1 \lambda_1 e^{\lambda_1 t} |w_1|) \frac{w_1}{|w_1|} + \text{sgn}(c_2 \lambda_2 e^{\lambda_2 t} |w_2|) \frac{w_2}{|w_2|} \right) \\ &= (k_1 \vec{e}_1 + k_2 \vec{e}_2) \cdot (\text{sgn}(k_1) \vec{e}_1 + \text{sgn}(k_2) \vec{e}_2) \quad (\text{where } k_1 \equiv c_1 \lambda_1 e^{\lambda_1 t} |w_1|, \vec{e}_1 \equiv \frac{w_1}{|w_1|}) \\ &= |k_1| + |k_2| + (|k_1| + |k_2|) \text{sgn}(k_1 k_2) \vec{e}_1 \cdot \vec{e}_2 = (|k_1| + |k_2|) (1 + \text{sgn}(k_1 k_2) \vec{e}_1 \cdot \vec{e}_2) > 0 \end{aligned}$$

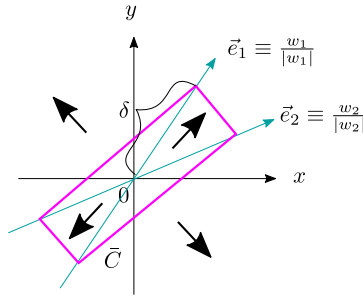


FIGURE 8. the simple closed curve  $\bar{\mathbf{C}}$  (Magenta) near  $(0,0)$  for  $\rho > 2$

(2)  $0 < \rho \leq 2, \Delta \leq 0$ , construct  $\bar{\mathbf{C}}$  as below  
it always holds  $(x', y')^T \cdot \vec{n} > 0$

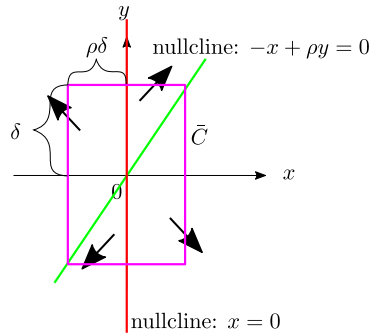


FIGURE 9. the simple closed curve  $\bar{\mathbf{C}}$  (Magenta) near  $(0,0)$  for  $0 < \rho \leq 2$

At each point of  $\bar{\mathbf{C}}$ , the field  $\mathbf{F} = (x', y')^T$  points toward the outside of  $\bar{\mathbf{C}}$

**Conclusion:** with **Poincare-Bendixson Ring Domain Theorem**,

there is a **closed trajectory (limit cycle)** inside  $\mathbf{R}$  between two simple closed curves  $\mathbf{C}$  and  $\bar{\mathbf{C}}$

JOURNAL.

(100-300 words, please type.) Give a description of the regular perturbation method in your own words. Discuss the idea behind the method, the purpose of the method and the limitations of the method.

**solution**

**Perturbation** develops an expression for the desired solution in terms of a formal power series in some "small" parameter  $\varepsilon$ , namely a perturbation series that quantifies the deviation from the exactly solvable problem.

$$y(\varepsilon, x_1, x_2, \dots) = \sum_{k=0}^{+\infty} y_k(x_1, x_2, \dots) \varepsilon^k$$

$y_0(x_1, x_2, \dots)$  is the known solution to the exactly solvable initial problem and  $y_1(x_1, x_2, \dots), y_2(x_1, x_2, \dots), \dots$  may be found iteratively by a mechanistic procedure. For small  $\varepsilon$  these higher-order terms in the series generally could become successively smaller.

Equations arising from mathematical models usually cannot be solved in exact form. The idea behind **perturbation** is that we breaks the problem into "solvable" and "perturbative" parts. Perturbation theory is widely used when the problem at hand does not have a known exact solution, but can be expressed as a "small" change to a known solvable problem. As a result, the computations of **perturbation** could be performed with a very high accuracy.

There are 2 limitations for regular perturbation

- The first one is called a **secular term**, like  $t \sin t$   
 In the approximation, the correction term  $\sum_{k=N}^{\infty} y_k(t) \varepsilon^k$  cannot be made arbitrarily small for  $t \in (0, +\infty)$  by choosing  $\varepsilon$  small enough.  
 The solution is the **Poincaré-Lindstedt** Method, that is the a scale transformation, to avoid the presence of secular terms in the expansion.
- The other one is that regular perturbation assumed a leading term of order unity, and it is not surprising that it missed some roots. The roots are different order, and one expansion does not reveal both.  
 The solution is **dominant balancing**, that means we examine each term carefully and determine which ones combine to give a dominant balance.