

BONUS.

16. (Circuits) An RCL circuit with a nonlinear resistance, where the voltage drop across the resistor is a nonlinear function of current, can be modeled by the Van der Pol equation

$$x'' + \rho(x^2 - 1)x' + x = 0$$

where ρ is a positive constant, and $x(t)$ is the current.

- In the phase plane, show that the origin is an unstable equilibrium.
- Sketch the nullclines and the vector field. What are the possible dynamics? Is there a limit cycle?

(Problem 16ab on Page 111, PDF Page 134)

solution

With the substitution $y = x'$

$$x' = y$$

$$y' = x'' = -x - \rho(x^2 - 1)y$$

The nullclines: $0 = y, 0 = -x - \rho(x^2 - 1)y$
the only critical point is $(0, 0)$

For the direction in the separated regions:

region 1: $y > 0, -x - \rho(x^2 - 1)y > 0$

$(x', y') = (y, -x - \rho(x^2 - 1)y) = (+, +)$

region 2: $y < 0, -x - \rho(x^2 - 1)y > 0$

$(x', y') = (y, -x - \rho(x^2 - 1)y) = (-, +)$

region 3: $y < 0, -x - \rho(x^2 - 1)y < 0$

$(x', y') = (y, -x - \rho(x^2 - 1)y) = (-, -)$

region 4: $y > 0, -x - \rho(x^2 - 1)y < 0$

$(x', y') = (y, -x - \rho(x^2 - 1)y) = (+, -)$

Consider the Jacobian J

$$J = \begin{pmatrix} \frac{\partial P}{\partial x} & \frac{\partial P}{\partial y} \\ \frac{\partial Q}{\partial x} & \frac{\partial Q}{\partial y} \end{pmatrix} = \begin{pmatrix} 0 & 1 \\ -1 - 2\rho xy & -\rho(x^2 - 1) \end{pmatrix}$$

a) At the critical point $(0, 0)$,

$$J = \begin{pmatrix} 0 & 1 \\ -1 & \rho \end{pmatrix}, \quad \det(J) \neq 0, \quad |J - \lambda I| = \lambda^2 - \rho\lambda + 1 = \lambda^2 - p\lambda + q = 0$$

With **Vieta theorem**, $\Delta = \rho^2 - 4, p = \lambda_1 + \lambda_2 = \text{Re}(\lambda_1) + \text{Re}(\lambda_2) = \rho > 0, q = 1$

if stable $\Rightarrow \text{Re}(\lambda_1) \leq 0, \text{Re}(\lambda_2) \leq 0 \Rightarrow \text{Re}(\lambda_1) + \text{Re}(\lambda_2) \leq 0$, there is a conflict, $(0, 0)$ is **unstable**

b) The the nullclines and the vector field are displayed in the figure above

(1) $\Delta = \rho^2 - 4 = 0 \Rightarrow \rho = 2$

repeated eigenvalues $\lambda_1 = \lambda_2 = 1$, solution $c_1 w e^{\lambda_1 t} + c_2(w + vt)e^{\lambda_1 t}$

the critical point $(0, 0)$ is an **unstable improper node**

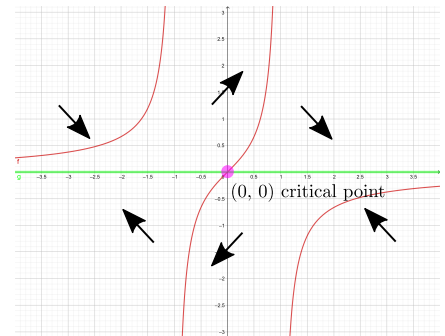


FIGURE 6. phase diagram $\rho = 1$

- (2) $\Delta = \rho^2 - 4 > 0 \Rightarrow 2 < \rho$
 real eigenvalues $\lambda_1 > \lambda_2 > 0$
 the critical point $(0, 0)$ is an **unstable node**
- (3) $\Delta = \rho^2 - 4 < 0 \Rightarrow 0 < \rho < 2$
 complex eigenvalues $\text{Re}(\lambda_1) = \text{Re}(\lambda_2) = \frac{\rho}{2} = \frac{\rho}{2} > 0$
 the critical point $(0, 0)$ is an **unstable spiral**

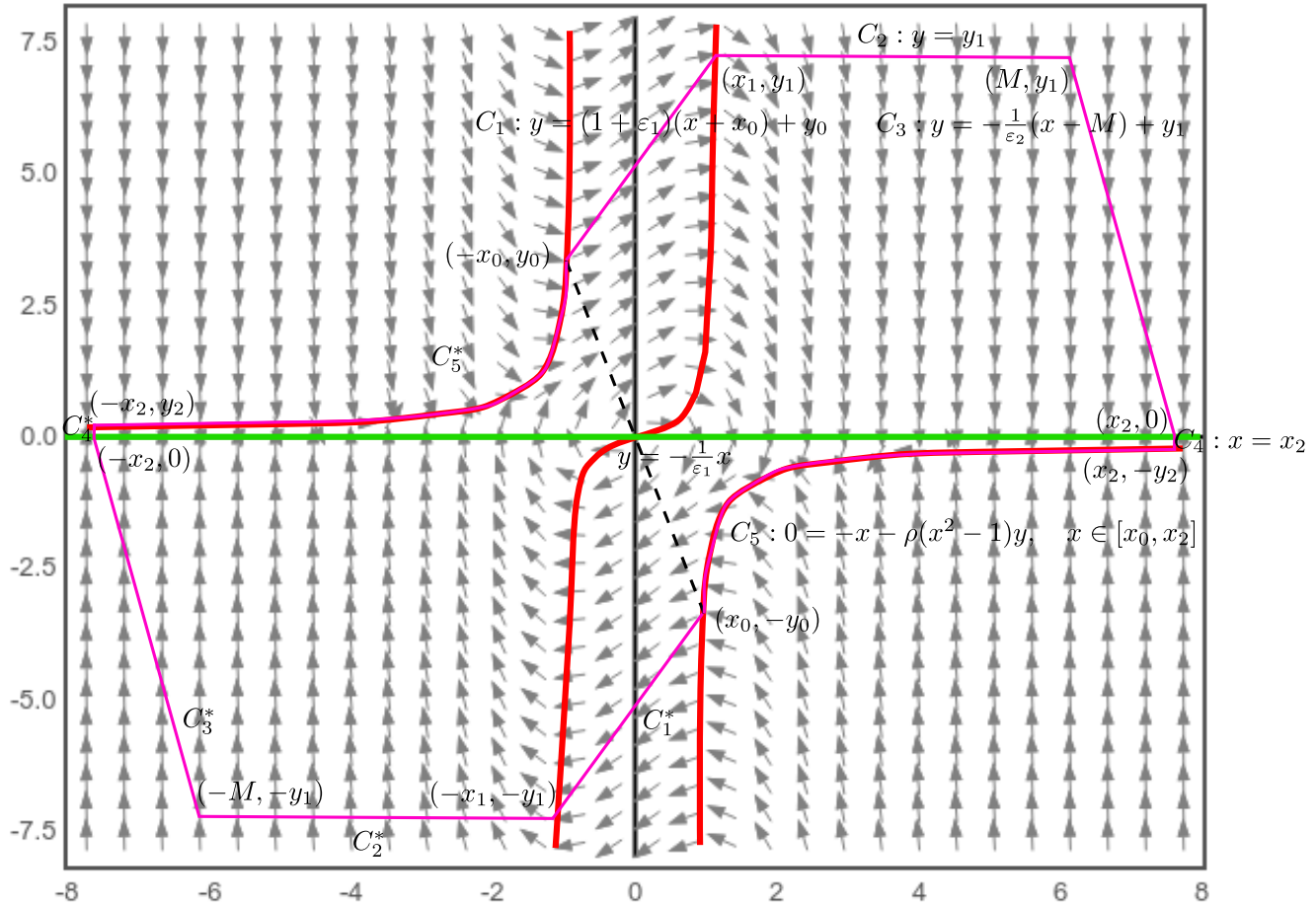


FIGURE 7. the simple closed curve C , where $(x', y')^T \cdot \vec{n} < 0$

(Note: **Green**: x nullcline; **Red**: y nullcline; **Magenta**: the simple closed curve C)

Theorem (Poincare-Bendixson Ring Domain Theorem). Suppose R is the finite region of the plane lying between two simple closed curves C and \bar{C} , and F is the velocity vector field for the system $x' = f(x, y)$ $y' = g(x, y)$. If

- (i) at each point of C and \bar{C} , the field F points toward the interior of R , and
- (ii) R contains no critical points,

then the system has a **closed trajectory** lying inside R

See the [MIT limit cycle note](https://math.mit.edu/~jorloff/supnotes/supnotes03/lc.pdf) or go to the url: <https://math.mit.edu/~jorloff/supnotes/supnotes03/lc.pdf>

For **Poincare-Bendixson Ring Domain Theorem**, we construct the simple closed curves first. curve \mathbf{C} consists of C_1, C_2, C_3, C_4, C_5 and the symmetrical curves $C_1^*, C_2^*, C_3^*, C_4^*, C_5^*$ of $(0, 0)$. For (x, y) on \mathbf{C} , always holds $(x', y')^T \cdot \vec{n} < 0$, here (x', y') : velocity field, \vec{n} : normal vector at (x, y) . Namely, for all (x, y) on \mathbf{C} , velocity field (x', y') points toward the inside of \mathbf{C} . Short explanation for $(x', y')^T \cdot \vec{n} < 0$ on C_1, C_2, C_3, C_4, C_5 as follows:

(1) $C_1 : y = (\rho + \varepsilon_1)(x + x_0) + y_0$

where $(-x_0, y_0)$ is the intersection of $y = -\frac{1}{\varepsilon_1}x$ and left branch of $0 = -x - \rho(x^2 - 1)y$

ε_1 must satisfy, at $(-x_0, y_0)$: $\frac{dy}{dx}|_{(-x_0, y_0)}$ slope of $0 = -x - \rho(x^2 - 1)y > (\rho + \varepsilon_1)$ slope of C_1

$$(-x_0, y_0) = \left(-\sqrt{1 + \frac{\varepsilon_1}{\rho}}, \sqrt{1 + \frac{\varepsilon_1}{\rho}}/\varepsilon_1 \right), \quad \text{left brach } 0 = -x - \rho(x^2 - 1)y \Leftrightarrow y = -\frac{1}{2\rho} \left(\frac{1}{x-1} + \frac{1}{x+1} \right)$$

$$\frac{dy}{dx} = \frac{1}{\rho} \frac{x^2 + 1}{(x^2 - 1)^2}, \quad \frac{dy}{dx}|_{(-x_0, y_0)} = \frac{1}{\rho} \frac{2 + \frac{\varepsilon_1}{\rho}}{\left(\frac{\varepsilon_1}{\rho}\right)^2} = \frac{2\rho + \varepsilon_1}{\varepsilon_1^2} > (\rho + \varepsilon_1) \Leftrightarrow \rho > 0 > -\varepsilon_1 \left(\frac{1 - \varepsilon_1^2}{2 - \varepsilon_1^2} \right)$$

We can find $0 < \varepsilon_1 < 1$ to satisfy it, for $\forall \rho > 0$, it always holds

$$(x', y')^T = (y, -x - \rho(x^2 - 1)y)^T, \quad \vec{n} = (-\rho - \varepsilon_1, 1)^T$$

$$(x', y')^T \cdot \vec{n} = -[(\varepsilon_1 y + x) + \rho x^2] < 0$$

(2) $C_2 : y = y_1$

where (x_1, y_1) is the intersection of C_1 and center branch of $0 = -x - \rho(x^2 - 1)y$

for (x, y) on C_2 , it has $x > x_1 > 0, y = y_1 > 0$, always holds

$$(x', y')^T = (y, -x - \rho(x^2 - 1)y)^T, \quad \vec{n} = (0, 1)^T$$

$$(x', y')^T \cdot \vec{n} = -x - \rho(x^2 - 1)y = [-x_1 - \rho(x_1^2 - 1)y_1] - (x - x_1) - (x^2 - x_1^2)y_1 = -(x - x_1) - (x^2 - x_1^2)y_1 < 0$$

(3) $C_3 : y = -\frac{1}{\varepsilon_2}(x - M) + y_1$

where (M, y_1) satisfies $M > \sqrt{1 + \frac{1}{\rho\varepsilon_2}}$

for (x, y) on C_3 , it has $x > M, y > 0$, always holds

$$(x', y')^T = (y, -x - \rho(x^2 - 1)y)^T, \quad \vec{n} = (1, \varepsilon_2)^T$$

$$(x', y')^T \cdot \vec{n} = -\varepsilon_2 x - y [\rho\varepsilon_2(x^2 - 1) - 1] = -\varepsilon_2 x - y [\rho\varepsilon_2(M^2 - 1) - 1] - y\rho\varepsilon_2(x^2 - M^2) < 0$$

(4) $C_4 : x = x_2$

where $(x_2, 0)$ is the intersection of C_3 and $y = 0$

for (x, y) on C_4 , it has $x = x_2, y < 0$, always holds

$$(x', y')^T = (y, -x - \rho(x^2 - 1)y)^T, \quad \vec{n} = (1, 0)^T$$

$$(x', y')^T \cdot \vec{n} = y < 0$$

(5) $C_5 : 0 = -x - \rho(x^2 - 1)y$ right branch

where $(x_2, -y_2)$ is the intersection of C_4 and $0 = -x - \rho(x^2 - 1)y$ right branch

for (x, y) on C_5 , it has $x > x_0 = \sqrt{1 + \frac{\varepsilon_1}{\rho}} > 1, y < 0$, always holds

$$y = -\frac{1}{2\rho} \left(\frac{1}{x-1} + \frac{1}{x+1} \right), \quad \frac{dy}{dx} = \frac{1}{\rho} \frac{x^2 + 1}{(x^2 - 1)^2}$$

$$(x', y')^T = (y, -x - \rho(x^2 - 1)y)^T = (y, 0)^T, \quad \vec{n} = \left(\frac{dy}{dx}, -1 \right)^T$$

$$(x', y')^T \cdot \vec{n} = - \left[\frac{1}{\rho} \frac{x^2 + 1}{(x^2 - 1)^2} \right] (-y) < 0$$

To sum up, $(x', y')^T \cdot \vec{n} < 0$ holds on C_1, C_2, C_3, C_4, C_5 , the velocity field is symmetrical to $(0, 0)$. For the symmetrical curves $C_1^*, C_2^*, C_3^*, C_4^*, C_5^*$ of $(0, 0)$, $(x', y')^T \cdot \vec{n} < 0$ holds on $C_1^*, C_2^*, C_3^*, C_4^*, C_5^*$. Thus, $(x', y')^T \cdot \vec{n} < 0$ holds for almost all points on entire \mathbf{C}

At each point of \mathbf{C} , the field $\mathbf{F} = (x', y')^T$ points toward the inside of \mathbf{C}

Now, consider to construct the simple closed curve $\bar{\mathbf{C}}$ (Magenta) near the critical point $(0, 0)$

(1) $\rho > 2, \Delta < 0$, construct $\bar{\mathbf{C}}$ as below

eigenvalues $\lambda_{1,2} = \frac{p \pm \sqrt{p^2 - 4}}{2} > 0$, with $(J - \lambda_{1,2})w_{1,2} = 0$, find $w_{1,2} = [1, \frac{p \pm \sqrt{p^2 - 4}}{2}]^T$
then $X(t) = (x, y)^T = c_1 e^{\lambda_1 t} w_1 + c_2 e^{\lambda_2 t} w_2$, $(x', y')^T = c_1 \lambda_1 e^{\lambda_1 t} w_1 + c_2 \lambda_2 e^{\lambda_2 t} w_2$
notice that coefficient $(c_1 \lambda_1 e^{\lambda_1 t}, c_2 \lambda_2 e^{\lambda_2 t})$ has the same signs as $(c_1 e^{\lambda_1 t}, c_2 e^{\lambda_2 t})$, (c_1, c_2)
the normal vector at (x, y) is $\vec{n} = \text{sgn}(c_1) \frac{w_1}{|w_1|} + \text{sgn}(c_2) \frac{w_2}{|w_2|}$, it always holds

$$\begin{aligned} (x', y')^T \cdot \vec{n} &= (c_1 \lambda_1 e^{\lambda_1 t} w_1 + c_2 \lambda_2 e^{\lambda_2 t} w_2) \cdot \left(\text{sgn}(c_1) \frac{w_1}{|w_1|} + \text{sgn}(c_2) \frac{w_2}{|w_2|} \right) \\ &= \left(c_1 \lambda_1 e^{\lambda_1 t} |w_1| \frac{w_1}{|w_1|} + c_2 \lambda_2 e^{\lambda_2 t} |w_2| \frac{w_2}{|w_2|} \right) \cdot \left(\text{sgn}(c_1 \lambda_1 e^{\lambda_1 t} |w_1|) \frac{w_1}{|w_1|} + \text{sgn}(c_2 \lambda_2 e^{\lambda_2 t} |w_2|) \frac{w_2}{|w_2|} \right) \\ &= (k_1 \vec{e}_1 + k_2 \vec{e}_2) \cdot (\text{sgn}(k_1) \vec{e}_1 + \text{sgn}(k_2) \vec{e}_2) \quad (\text{where } k_1 \equiv c_1 \lambda_1 e^{\lambda_1 t} |w_1|, \vec{e}_1 \equiv \frac{w_1}{|w_1|}) \\ &= |k_1| + |k_2| + (|k_1| + |k_2|) \text{sgn}(k_1 k_2) \vec{e}_1 \cdot \vec{e}_2 = (|k_1| + |k_2|) (1 + \text{sgn}(k_1 k_2) \vec{e}_1 \cdot \vec{e}_2) > 0 \end{aligned}$$

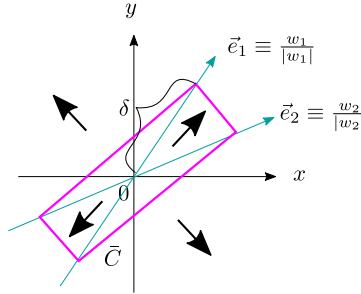


FIGURE 8. the simple closed curve $\bar{\mathbf{C}}$ (Magenta) near $(0,0)$ for $\rho > 2$

(2) $0 < \rho \leq 2, \Delta \leq 0$, construct $\bar{\mathbf{C}}$ as below
it always holds $(x', y')^T \cdot \vec{n} > 0$

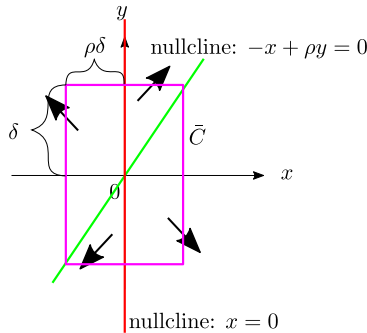


FIGURE 9. the simple closed curve $\bar{\mathbf{C}}$ (Magenta) near $(0,0)$ for $0 < \rho \leq 2$

At each point of $\bar{\mathbf{C}}$, the field $\mathbf{F} = (x', y')^T$ points toward the outside of $\bar{\mathbf{C}}$

Conclusion: with **Poincare-Bendixson Ring Domain Theorem**,

there is a **closed trajectory (limit cycle)** inside \mathbf{R} between two simple closed curves \mathbf{C} and $\bar{\mathbf{C}}$